"Measurement of the Differential Cross Section dσ/d(cos θt) for Top-Quark Pair Production in p-pbar Collisions at sqrt(s) = 1.96 TeV"

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Abstract
We report a measurement of the differential cross section, d(σ)/d(cos θt), for top-quark-pair production as a function of the top-quark production angle in proton-antiproton collisions at sqrt(s) = 1.96 TeV. This measurement is performed using data collected with the CDF II detector at the Tevatron, corresponding to an integrated luminosity of 9.4/ fb. We employ the Legendre polynomials to characterize the shape of the differential cross section at the parton level. The observed Legendre coefficients are in good agreement with the prediction of the next-to-leading-order standard-model calculation, with the exception of an excess linear-term coefficient, a1 = 0.40 ± 0.12, compared to the standard-model prediction of a1 = 0.15^{+0.07}_{-0.03}.

Document type : Article de périodique (Journal article)

Référence bibliographique


DOI : 10.1103/PhysRevLett.111.182002
Measurement of the Differential Cross Section \(d\sigma/d(\cos\theta_f)\) for Top-Quark Pair Production in \(p\bar{p}\) Collisions at \(\sqrt{s} = 1.96\) TeV

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We report a measurement of the differential cross section $d\sigma/d(\cos \theta)$ for top-quark pair production as a function of the top-quark production angle in proton-antiproton collisions at $\sqrt{s} = 1.96$ TeV. This measurement is performed using data collected with the CDF II detector at the Tevatron, corresponding to an integrated luminosity of $9.4 \text{ fb}^{-1}$. We employ the Legendre polynomials to characterize the shape of the differential cross section at the parton level. The observed Legendre coefficients are in good agreement with the prediction of the next-to-leading-order standard-model calculation, with the exception of an excess linear-term coefficient $a_1 = 0.40 \pm 0.12$, compared to the standard-model prediction of $a_1 = 0.15^{+0.07}_{-0.03}$.

The Collider Detector at Fermilab (CDF) and D0 experiments have measured an anomalously large forward-backward asymmetry ($A_{FB}$) in top-quark pair ($t\bar{t}$) hadroproduction. The latest measurements are $A_{FB} = (16.4 \pm 4.5)\%$ from CDF [1] and $A_{FB} = (19.6 \pm 6.5)\%$ from D0 [2]. This asymmetry is the manifestation of a charge asymmetry in $t\bar{t}$ production via the CP-even [3] initial state at the Fermilab Tevatron proton-antiproton ($p\bar{p}$) collider. The standard model (SM) [4] predicts a small forward-backward asymmetry ($8.8 \pm 0.6)\%$ at next-to-leading order (NLO) in the strong coupling constant $\alpha_s$ [5,6]. The tension between the Tevatron measurements and the predictions has stimulated new work on the SM calculation [7–10] and on possible non-SM sources for the asymmetry [11]. The charge asymmetry is also under study at the LHC, but any effects are expected to be much smaller due to the forward-backward symmetric (proton-proton) initial state [12], and the results are so far inconclusive [13,14].

We measure the differential cross section $d\sigma/d(\cos \theta)$, where $\theta$ is the angle between the top-quark momentum and the incoming proton momentum, as measured in the $t\bar{t}$ center-of-mass frame. The inclusive measurements of $A_{FB}$ are equivalent to a two-bin measurement of this differential cross section, with one bin forward ($\cos \theta > 0$) and one bin backward ($\cos \theta < 0$). The full shape of the differential cross section provides additional information and has the potential to discriminate among various calculations of the SM as well as models of non-SM physics. One of the aims of this study is to identify what aspects of the shape of $d\sigma/d(\cos \theta)$ explain the $A_{FB}$.

We characterize the shape of $d\sigma/d(\cos \theta)$ by employing the Legendre polynomials [15], which are fundamental to the general theory of scattering of particles in the present spin-averaged case [16]. The orthonormality of these polynomials on the interval $[-1, 1]$ allows a unique decomposition of the cross section into a Legendre polynomial series. We write

$$\frac{d\sigma}{d(\cos \theta)} = \sum_{\ell=0}^{\infty} a_{\ell} P_{\ell}(\cos \theta),$$

(1)

where $P_{\ell}$ is the Legendre polynomial of degree $\ell$, and $a_{\ell}$ is the Legendre moment of degree $\ell$. Because the experimental sensitivity degrades as $\ell$ increases, we restrict the sum to $\ell \leq 8$. Since the moment $a_0$ contains only the total cross section, we scale all the moments so that $a_0 = 1$.

At leading order (LO) in the SM, the differential cross section for $q\bar{q} \rightarrow t\bar{t}$ is

$$\frac{d\sigma_{q\bar{q} \rightarrow t\bar{t}}}{d\Omega}(\cos \theta, \hat{s}) = \frac{\beta \alpha_s^3}{144 \pi \hat{s}} [2 - \beta^2(1 + \cos^2 \theta)],$$

(2)
where $\beta$ is the velocity of the top quark in units of $c$ [17] and $s$ is the Mandelstam variable [18]. After integrating over $s$ to obtain $d\sigma/d(\cos\theta_t)$ and comparing to the Legendre polynomials, we expect nonzero values only for $a_0$ and $a_2$. The addition of the $gg \rightarrow t\bar{t}$ process is expected to add small contributions to all the even-degree Legendre moments. We study the LO SM via a sample of simulated events generated by PYTHIA [19]. At next-to-leading order in the SM, additional contributions to all the Legendre moments appear, including the odd moments. These nonzero odd moments introduce the lowest-order contributions to $A_{FB}$. The NLO SM theoretical calculation adopted in this Letter includes the full effects of both quantum chromodynamics and the electroweak theory [5].

A wide variety of non-SM proposals has been put forward to explain the large value of $A_{FB}$ observed at the Tevatron. These form two broad classes, depending on whether the new physics is dominated by $s$- or $t$-channel exchange. In order to characterize the effect of these models on the differential cross section, we study two representative models. An $s$-channel model “Octet $A$” hypothesizes the existence of a heavy ($m_{Z'} = 2$ TeV/$c^2$) partner of the gluon with axial-vector couplings to quarks [20]. This produces an enhanced linear-term coefficient $a_1$ in $d\sigma/d(\cos\theta_t)$ [17]. A $t$-channel model “$Z'$ 200” contains a new, heavy ($m_{Z'} = 200$ GeV/$c^2$) vector boson with a flavor changing $u$-$Z'$-$t$ coupling [21]. The resulting additional term in the cross section has a leading dependence $s/t = 1/(1 - \cos\theta_t)$, where $t$ is the Mandelstam variable [18]. This behavior produces large Legendre moments at all degrees. These leading behaviors are generic predictions of $s$- and $t$-channel models [17]. Both models are studied via samples of simulated events generated at LO by MADGRAPH [22]. The LO and NLO SM calculations, as well as these two benchmark non-SM models, are shown in Figs. 1 and 2.

We study the full sample of top-quark pair candidate events in the decay channel with a single lepton in the final state collected by the CDF experiment during Run II of the Fermilab Tevatron. The CDF II detector is a general purpose particle detector employing a large charged-particle tracking volume inside a solenoidal magnetic field coaxial with the beam direction, surrounded by calorimeters and muon detectors [23,24]. The collected data correspond to an integrated luminosity of 9.4 fb$^{-1}$ of $p\bar{p}$ collisions. The general features of the event selection requirements are as follows. We require exactly one well-reconstructed charged-lepton candidate (electron or muon) with $p_T > 20$ GeV/$c$, an imbalance in the total event transverse momentum (missing transverse energy [25]) $E_T > 20$ GeV, and four or more calorimeter-energy clusters (jets [26]), three with $E_T > 20$ GeV and the fourth with $E_T > 12$ GeV, in the central part of the detector ($|\eta| < 2.0$). We further require that at least one of the jets be identified (tagged) as having a displaced vertex resulting from the decay of a bottom-quark meson, which is

![FIG. 1 (color online). The predicted differential cross sections of the LO SM [19], the NLO SM [5], and the benchmark models for $s$- and $t$-channel new physics [20,21]. The band around the NLO SM prediction represents the uncertainty due to renormalization-scale choice.](image1)

produced from the dominant top-quark decay $t \rightarrow Wb$. Further details on the on-line and off-line event selection requirements are in Ref. [1]. The resulting data set is enriched in $t\bar{t}$ events, but it contains non-$t\bar{t}$ background events as well, dominated by events in which a $W$ boson is produced in association with hadron jets. The rates and differential distributions of all the sources of non-$t\bar{t}$ backgrounds are well understood [1]. We expect to observe $2750 \pm 427$ $t\bar{t}$ events and $1026 \pm 210$ non-$t\bar{t}$ background events, and we observe 3864 $t\bar{t}$ candidate events.

We reconstruct the top quark and the top antiquark from their decay products, using the measured momentum of the lepton and the four jets, as well as the missing transverse energy. We fit each possible jet-to-parton assignment to the $t\bar{t}$ hypothesis. We require that two of the jets be consistent with the decay of a $W$ boson of mass 80.4 GeV/$c^2$ and that

![FIG. 2 (color online). Measured Legendre moments $a_1$–$a_4$, with theory predictions overlaid.](image2)
the lepton and missing transverse energy also be consistent with the decay of a $W$ boson. We further require that each reconstructed $W$ boson, when paired with one of the remaining jets, be consistent with the decay of a top quark of mass 172.5 GeV/$c^2$ [27]. The jet-to-parton assignment which is most consistent with this $t\bar{t}$ hypothesis is used to calculate the top-quark production angle as measured in the detector $\cos\theta_t^{\text{det}}$ for each event.

We exploit the orthonormality of the Legendre polynomials to estimate the Legendre moments without performing a fit. Given a distribution $f(\cos\theta)$, the Legendre moments of $f$ are

$$a_\ell = \frac{2\ell + 1}{2} \int_{-1}^{1} d(\cos\theta)f(\cos\theta)P_\ell(\cos\theta).$$

(3)

The data are described by an empirical distribution [28] $f(\cos\theta^{\text{det}}) = \sum \delta(\cos\theta^{\text{det}} - \cos\theta_i^{\text{det}})$, where $\delta(x)$ is the Dirac $\delta$ function and the index $i$ runs over the events in the data set. Using this distribution in Eq. (3) greatly simplifies the integration due to the Dirac delta functions, so the moments of the observed $\cos\theta_t^{\text{det}}$ distribution are

$$a_\ell^{\text{det}} = \frac{2\ell + 1}{2} \sum P_\ell(\cos\theta^{\text{det}}).$$

(4)

Then, the estimate of the moments is

$$a_\ell = \sum K_{\ell m} \left( \sum P_m(\cos\theta^{\text{det}}) - a_m^{BG} \right),$$

(5)

where $a_m^{BG}$ represents the Legendre moments of the distribution of $\cos\theta^{\text{det}}$ predicted by the background model, and $K_{\ell m}$ is a correction matrix that accounts for the finite resolution of the detector and for the nonuniform detector acceptance and selection efficiency. The matrix $K$ is developed from a sample of fully simulated $t\bar{t}$ Monte Carlo events generated by the POWHEG NLO SM generator [29]. It describes the response of the detector and the effects of the event selection requirements. No smoothing or regularization is applied in this correction procedure, in contrast to the correction procedure of Ref. [1].

The statistical uncertainties on the moments are given by a root-mean-square covariance matrix including correlations. In order to estimate the effect from each of several sources of systematic uncertainty in the model assumptions, we vary the corresponding nuisance parameter that alters either the background prediction or correction matrix, and then perform the full correction procedure again. The resulting parton-level moments estimate is compared to the unvaried moments, and then the covariance matrix describing the uncertainty on the measurement is $\sigma_{\ell m} = \delta_m \delta_\ell$, where $\delta_\ell = a_\ell^{\text{varied}} - a_\ell^{\text{nominal}}$. We study systematic shifts due to the uncertainty in the jet-energy scale, the rate of the backgrounds, the shape of the backgrounds, the modeling of parton showering, the modeling of color reconnection, the modeling of initial- and final-state radiation, and the parton distribution functions of the proton and antiproton. We sum the resulting covariance matrices and add them to the statistical covariance matrix to obtain a covariance matrix that fully describes the uncertainty of the measurement of the parton-level Legendre moments. The eigenvalues and eigenvectors of the covariance matrix [30] can be used to calculate a $\chi^2$ goodness-of-fit statistic with 8 degrees of freedom in order to perform fits to the data.

The parton-level Legendre moments are shown in Fig. 2 and in Table I. We observe good agreement within the uncertainties with the NLO SM prediction for moments $a_2-a_8$, but $a_1$ is in excess of the prediction. That is, a mild excess is observed in the differential cross section in the term linear in $\cos\theta$, while all other terms are as predicted by the SM. The LO SM prediction is strongly disfavored by the linear term, with a significance of more than 3 standard deviations. The benchmark $t$-channel model $Z'$ 200 is disfavored by $a_2$ and $a_3$. The benchmark $s$-channel model Octet $A$ is in good agreement with the data.

We determine the contribution of each Legendre moment to the $A_{FB}$ from the inherent asymmetry of each polynomial (Fig. 3). The observed $A_{FB}$ (19.9 ± 5.7)% is completely dominated by the excess linear term $a_1 \cos\theta$, which contributes (20.1 ± 6.1)%). The $A_{FB}$ contributed by the nonlinear asymmetric terms $a_3$, $a_5$, and $a_7$ is negligible ($-0.2 ± 3.1)% and is consistent with the SM prediction (7.3% from the linear term, $-0.3%$ from the nonlinear terms). The correlation between the measurements of $A_{FB}$ from the linear and nonlinear terms is $-29%$.

A more traditional picture of the differential cross section (Fig. 4) is obtained by integrating the Legendre series over intervals (bins) in $\cos\theta$. This shows the fraction of the total cross section that accrues in each bin. The uncertainties are strongly correlated, and they are dominated by the large uncertainties on the high-degree Legendre moments.

Because the nonlinear moments $a_2$–$a_8$ are in good agreement with the uncertainties with the NLO SM prediction, we may obtain a more precise, but model-dependent, estimate of the linear term by explicitly

<table>
<thead>
<tr>
<th>$\ell$</th>
<th>$a_\ell$ (obs)</th>
<th>$a_\ell$ (pred)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>0.40 ± 0.12</td>
<td>0.15 ± 0.07</td>
</tr>
<tr>
<td>2</td>
<td>0.44 ± 0.25</td>
<td>0.28 ± 0.05</td>
</tr>
<tr>
<td>3</td>
<td>0.11 ± 0.21</td>
<td>0.030 ± 0.014</td>
</tr>
<tr>
<td>4</td>
<td>0.22 ± 0.28</td>
<td>0.035 ± 0.008</td>
</tr>
<tr>
<td>5</td>
<td>0.11 ± 0.33</td>
<td>0.005 ± 0.002</td>
</tr>
<tr>
<td>6</td>
<td>0.24 ± 0.40</td>
<td>0.006 ± 0.003</td>
</tr>
<tr>
<td>7</td>
<td>−0.15 ± 0.48</td>
<td>−0.003 ± 0.001</td>
</tr>
<tr>
<td>8</td>
<td>0.16 ± 0.65</td>
<td>−0.0019 ± 0.0003</td>
</tr>
</tbody>
</table>
assuming that the nonlinear moments are as predicted by
the NLO SM calculation. Using the covariance matrix and
the fitting procedure described in Ref. [30], we fit to the
measured moments, taking the NLO SM prediction for the
nonlinear moments with their scale uncertainties as a prior
assumption, obtaining
\[ a_1 = 0.39 \pm 0.11 \] (including statistical and systematic uncertainty). Through the correlations
among the measured moments, this reduces the uncertainty
on \( a_1 \) by about 10\% while shifting the central value less
than 3\%. The resulting curve is also shown in Fig. 4.

In conclusion, we have presented the first measurement
of the top-quark pair production differential cross section
\( d\sigma/d(\cos\theta_t) \) in \( pp \) collisions at \( \sqrt{s} = 1.96 \) TeV as a function
of the production angle of the top quark. In order to
probe the origin of the top-quark production asymmetry,
we decompose the angular form into Legendre
polynomials. We observe that the coefficient of the \( \cos\theta_t \)
term in the differential cross section \( a_1 = 0.40 \pm 0.12 \) is in
excess of the NLO SM prediction \( 0.15 \pm 0.003 \), while the
remainder of the differential cross section is in good agree-
ment within the uncertainties with the NLO SM prediction.
The top-quark forward-backward asymmetry is thus com-
pletely dominated by the linear term. The result constrains
\( t \)-channel explanations of the asymmetry and favors asym-
metry models with strong \( s \)-channel components.

We thank T. Tait, S. Jung, W. Bernreuther, and Z.-G. Si for their assistance in preparing the theoretical models and calculations used in this Letter, and T. Rizzo for helpful conversations. We also thank the development teams of SCIPY, PYTABLES, MATPLOTLIB, and IPYTHON for their useful tools [31]. We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science, and Technology of Japan; the Natural Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the Swiss National Science Foundation; the A. P. Sloan Foundation; the Bundesministerium für Bildung und Forschung, Germany; the Korean World Class University Program and the National Research Foundation of Korea; the Science and Technology Facilities Council and the Royal Society, U.K.; the Russian Foundation for Basic Research; the Ministerio de Ciencia e Innovación and Programa Consolider-Ingenio 2010, Spain; the Slovak R&D Agency; the Academy of Finland; the Australian Research Council (ARC); and the EU Community Marie Curie Fellowship Contract No. 302103.
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[3] CP is the simultaneous transformation of charge conjugation and parity.
[15] We normalize the Legendre polynomials so that \( \int_{-1}^{1} P_i(x)P_j(x)dx = \delta_{ij} \), where \( \delta_{ij} \) is the Kronecker delta.
[24] We use a cylindrical coordinate system with the origin at the center of the CDF detector, \( z \) pointing in the direction of the proton beam, \( \theta \) and \( \phi \) representing the polar and azimuthal angles, respectively, and pseudorapidity defined by \( \eta = -\ln \tan(\theta/2) \). The transverse momentum \( p_T \) (transverse energy \( E_T \)) is defined to be \( p \sin \theta (E \sin \theta) \).
[25] The calorimeter missing energy \( E_T^{\text{cal}} \) is defined by the sum over calorimeter cells \( E_T^{\text{cal}} = \sum_i E_T^i h_i \), where \( i \) is a calorimeter cell number with \( |\eta| < 3.3 \), and \( h_i \) is a unit vector perpendicular to the beam axis and pointing at the \( i \)th calorimeter cell. The reconstructed missing energy \( \vec{E}_T \) is derived by subtracting from \( \vec{E}_T \) (cal) the energies associated with components of the event not registered by the calorimeter, such as muons and jet-energy adjustments. \( \vec{E}_T \) (cal) and \( \vec{E}_T \) are the scalar magnitudes of \( \vec{E}_T \) (cal) and \( \vec{E}_T \), respectively.